

# New level information on $Z = 51$ isotopes, $^{111}\text{Sb}_{60}$ and $^{134,135}\text{Sb}_{83,84}$

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**Abstract.** New data for low-spin low-energy levels in  $^{111,134,135}\text{Sb}$  are presented. The observed structures are compared to shell-model calculations. The monopole shifts for the  $d_{5/2}$  and  $g_{7/2}$  single-proton levels and the spin-orbit splitting for the  $d_{5/2}$  and  $d_{3/2}$  orbitals as  $N/Z$  moves from  $\sim 1$  up to 1.6 are discussed.

**PACS.** 23.40.-s  $\beta$  decay; double  $\beta$  decay; electron and muon capture – 21.10.-k Properties of nuclei; nuclear energy levels – 21.60.-n Nuclear structure models and methods

## 1 Introduction

The systematic changes in structures of the low-spin states below 2.0 MeV of odd- $A$  Sb nuclides from  $^{101}\text{Sb}$  to  $^{135}\text{Sb}$  provide important information regarding the monopole shift of the  $g_{7/2}$  single-proton state between the  $N = 50$  and  $N = 82$  closed neutron shells and beyond. Levels below 2.0 MeV arise from coupling the  $2^+$  phonon of the Sn core to single-particle  $g_{7/2}$  and  $d_{5/2}$  states, and the presence or absence of degeneracy with the position of the  $2^+$  core level is one indication of the extent of configuration mixing.

For odd- $A$  Sb nuclei with  $115 \leq A \leq 133$ , the structures of states below 2.0 MeV have been established by a number of different experiments. Whereas, only a recent study of the structure of  $^{109}\text{Sb}$  by Ressler *et al.* [1], has provided new data for low-spin levels below  $^{113}\text{Sb}$ . In this paper new data for levels in  $^{111}\text{Sb}$  populated in the decay of  $^{111}\text{Te}$  are presented.

Above the  $N = 82$  shell closure, single-particle states have been identified in  $^{133}\text{Sb}$  [2], and the  $d_{5/2}$  single-particle state has been identified in  $^{135}\text{Sb}$  [3]. The positions of these single-particle levels are important for understanding the evolution of nuclear structure as  $N/Z$  ex-

ceeds 1.5 and the neutron dripline is approached. Also presented in this paper are new data for the low-spin states of both  $^{134}\text{Sb}$  and  $^{135}\text{Sb}$ .

## 2 Experimental

Low-spin states in  $^{111,113,115}\text{Sb}$  were populated via  $\beta^+/\text{EC}$  decay of  $^{111,113,115}\text{Te}$  nuclei that were produced at Argonne National Laboratory using the  $^{58}\text{Ni}(^{56}\text{Fe},2\text{pn})^{111}\text{Te}$ ,  $^{60}\text{Ni}(^{56}\text{Fe},2\text{pn})^{113}\text{Te}$ , and  $^{62}\text{Ni}(^{56}\text{Fe},2\text{pn})^{115}\text{Te}$  fusion-evaporation reactions. Reaction products were separated in the Fragment Mass Analyzer on the basis of their mass to charge ( $A/Q$ ) ratio, and following mass separation, the recoils were implanted in the tape of a moving tape collector (MTC). As the half-life of  $^{111}\text{Te}$  ( $T_{1/2} = 26.2(6)$  s) is shorter than nuclei produced from other reaction channels, the tape was moved periodically to a Pb-shielded counting station to maximize coincidence events; and to reduce contribution to the  $\gamma$  spectra from both the decays of daughter and granddaughter nuclides, and the decays of nuclides that are collected on the tape owing to similar  $A/Q$  values.

Levels in  $^{134,135}\text{Sb}$  were populated by  $\beta\text{dn}$  and  $\beta^-$  decay of  $^{135}\text{Sn}$ , respectively. Neutron-rich Sn nuclei were produced at ISOLDE, CERN by using a 1.4 GeV proton beam pulse to induce fission of a  $\text{UC}_2$  target. Sn isotopes were then selectively ionized using the Resonance Ionization Laser Ion Source which consists of three copper vapor

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pumped dye lasers tuned to two resonant atomic excitations and a third to an energy which allowed for ionization of the electron into the continuum. Following mass separation,  $^{135}\text{Sn}$  was implanted into the tape of a moving tape collector, data were collected, and the tape was moved to remove daughter, granddaughter, and isobaric nuclides prior to the next proton pulse. Data were collected with the “laser-on” and with the “laser-off” to distinguish which peaks in the  $\gamma$  spectra could be assigned to the decay of Sn, and which were a result of the decay of surface ionized  $^{135}\text{Cs}$ . In both experiments,  $\gamma$  singles,  $\gamma$ -time, and  $\gamma$ - $\gamma$  coincidences were collected.

### 3 Levels in proton-rich Sb nuclides

On the proton-rich side, eleven new states were identified in  $^{111}\text{Sb}$ , including tentative assignments of the yrast  $1/2^+$  and  $3/2^+$  levels at 487 and 881 keV, respectively [4]. These data were combined with similar levels in  $^{109}\text{Sb}$  [1] to improve the estimate for the positions of the  $s_{1/2}$  and  $d_{3/2}$  single-proton basis states in  $^{101}\text{Sb}$  that underlie nuclear structure calculations in this mass region. With the position of the  $d_{3/2}$  level at 2.9 MeV, the spin-orbit splitting in  $^{101}\text{Sb}$  can be seen to be about twice the 1477 keV separation known for the  $d_{5/2}$  to  $d_{3/2}$  separation in  $^{133}\text{Sb}$ , in agreement with the conjecture made by Schiffer *et al.* [5]. There are, of course, numerous other approaches to a full description of the spin-orbit splitting, however, we note that a further reduction of the  $d_{5/2}$  to  $d_{3/2}$  separation is found going from  $^{133}\text{Sb}_{82}$  to  $^{207}\text{Tl}_{126}$  where that separation is reduced to  $> 1332$  keV [6]. We include the “greater than” symbol to recognize that, although the  $3/2^+$  level at 351 keV has a large spectroscopic factor, the  $5/2^+$  level at 1653 keV does not contain all of the  $d_{5/2}$  strength. Hence, the centroid of the spin-orbit splitting is surely larger than 1332 keV, perhaps even approaching the 1477 keV splitting observed in  $^{133}\text{Sb}$ .

### 4 Levels in neutron-rich Sb nuclides

On the neutron-rich side, new level structures have been identified in  $^{135}\text{Sb}$  and  $^{134}\text{Sb}$  where the levels were populated following direct  $\beta^-$  and  $\beta$ -delayed neutron decays of  $^{135}\text{Sn}$ , respectively. The observed levels in  $^{135}\text{Sb}$  below 1.0 MeV are shown in fig. 1, along with the results of a shell model calculation using the CD Bonn interaction, where the single-proton  $d_{5/2}$  and  $d_{3/2}$  levels have been lowered by 300 keV as described previously, thereby holding the spin-orbit splitting the same as for  $^{133}\text{Sb}$  [3]. As can be seen, the positions for the calculated levels are in good agreement with the positions of the observed levels. The new levels populated in  $\beta$ -delayed neutron decay in odd-odd  $^{134}\text{Sb}$  below 1.5 MeV are also shown in fig. 1, along with the levels calculated in an OXBASH calculation using the KH5082 interaction. Levels in  $^{134}\text{Sb}$  at 13, 330, and 383 keV were previously reported by Korgul *et al.* [7]. Again, an excellent fit is found, in spite of the fact that the KH5082 interaction for this mass region has been imported from the  $^{208}\text{Pb}$  region and scaled by  $A^{-1/3}$ .

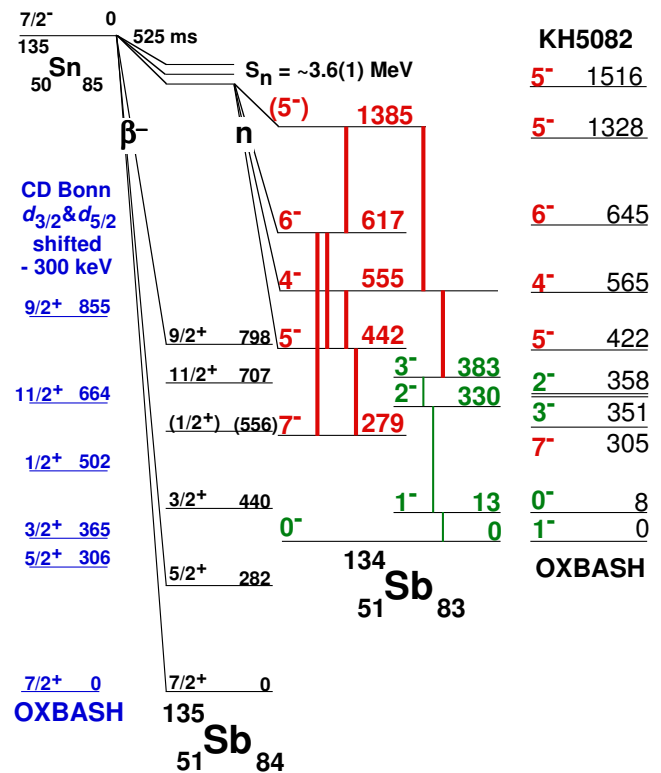


Fig. 1. The  $\beta^-$  and  $\beta$ dn decays of  $^{135}\text{Sn}$ . The structures of  $^{134}\text{Sb}$  and  $^{135}\text{Sb}$  are compared with shell model calculations with the parameters described in the text.

### 5 Conclusions

These new data have provided improved values for the monopole shifts of single-proton states from  $^{109}\text{Sb}_{58}$  through  $^{135}\text{Sb}_{85}$ . For both light and heavy Sb nuclides, shell model calculations provide a good description of the observed structures. It should be noted that other approaches to monopole shifts in neutron-rich nuclides have recently been published by both Otsuka *et al.* [8], and by Hamamoto [9].

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